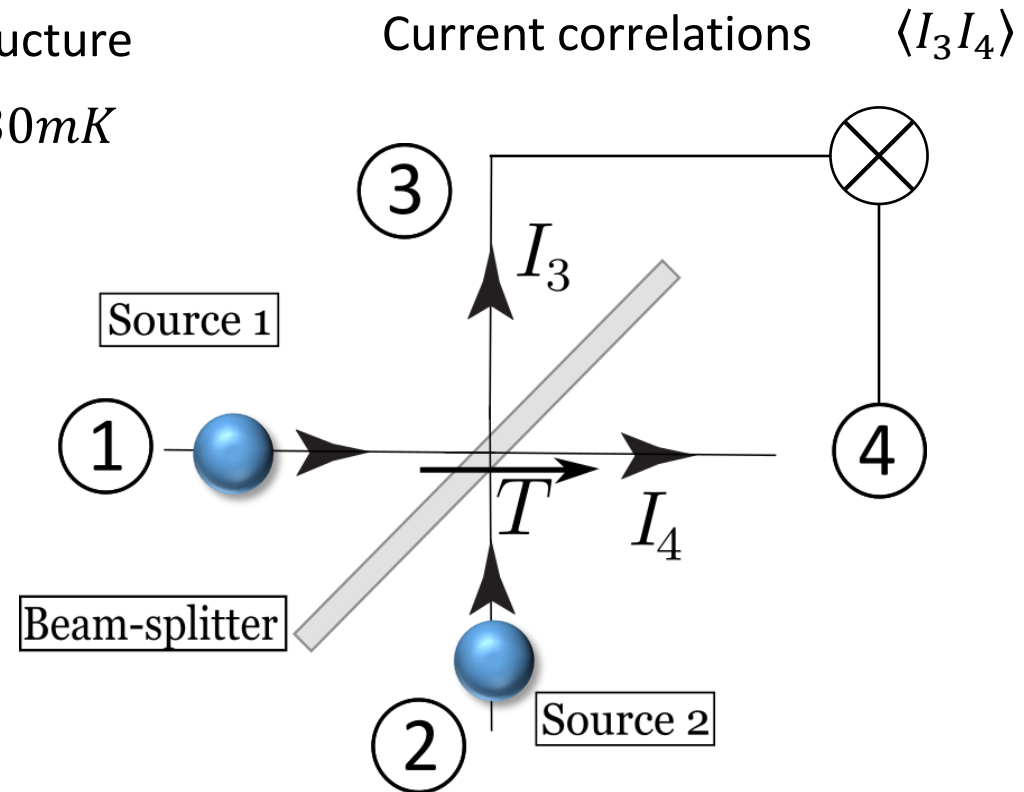
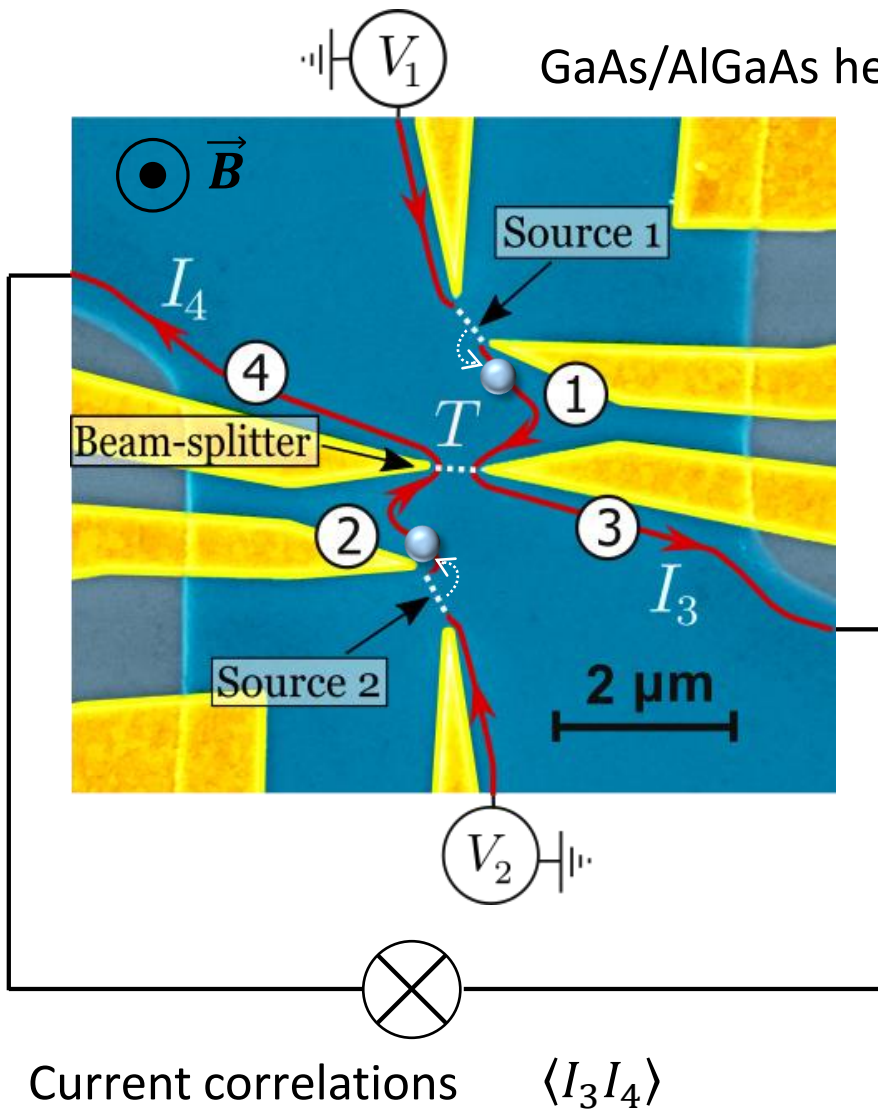


G. Fève, Laboratoire de Physique de l'ENS

Electron optics experiments in quantum Hall conductors

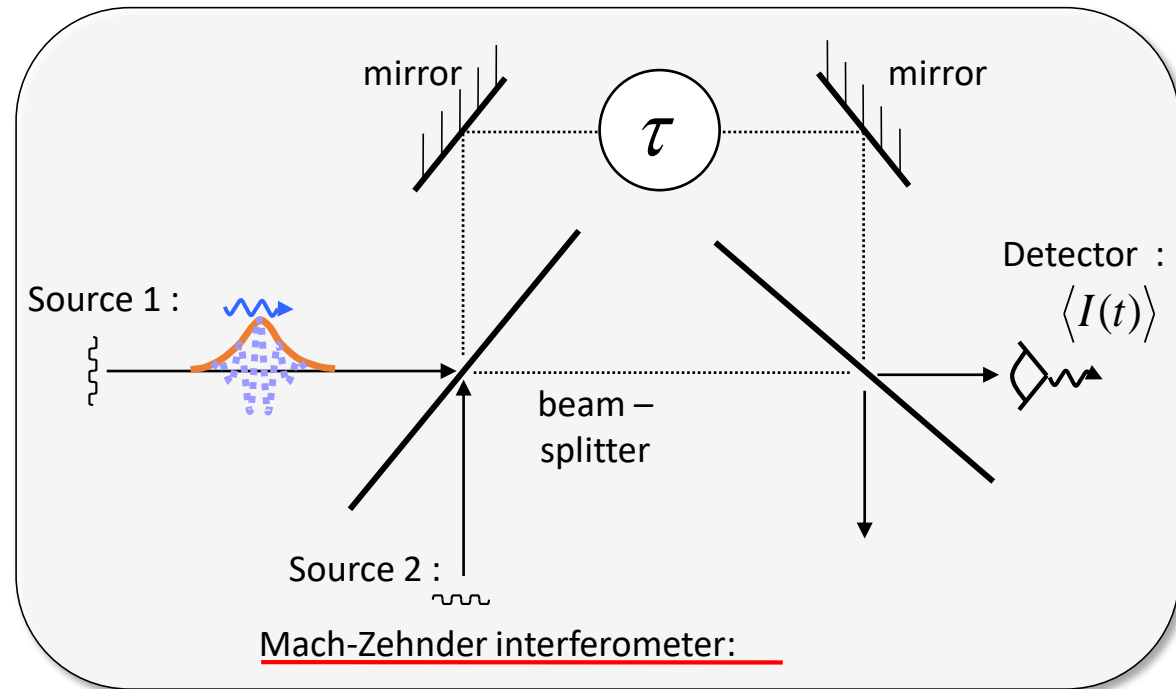


Optics: $\begin{cases} E(t) & \text{Electric field} \\ I(t) & \text{light intensity} \end{cases}$

Single particle interferometer

$$G^{(1)}(t + \tau, t) \propto \langle E(t + \tau)E(t) \rangle$$

Coherence of electric field

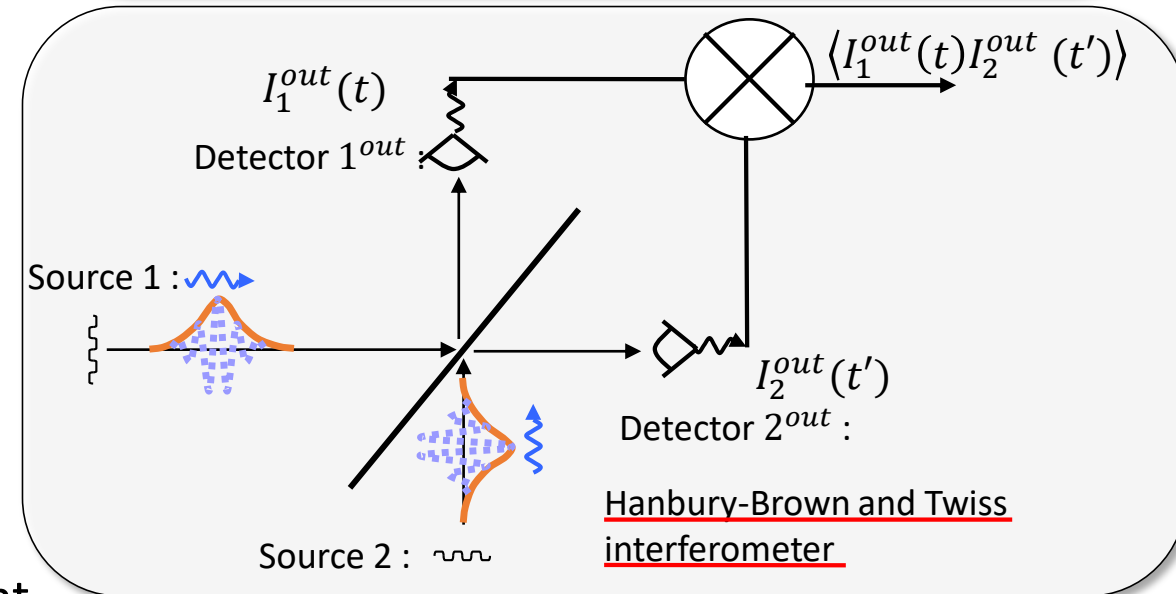


Two-particle interferometer

$$\langle I_1^{out}(t) I_2^{out}(t') \rangle$$

$$\propto \langle E_1^{in}(t) E_1^{in}(t') \rangle \langle E_2^{in}(t) E_2^{in}(t') \rangle$$

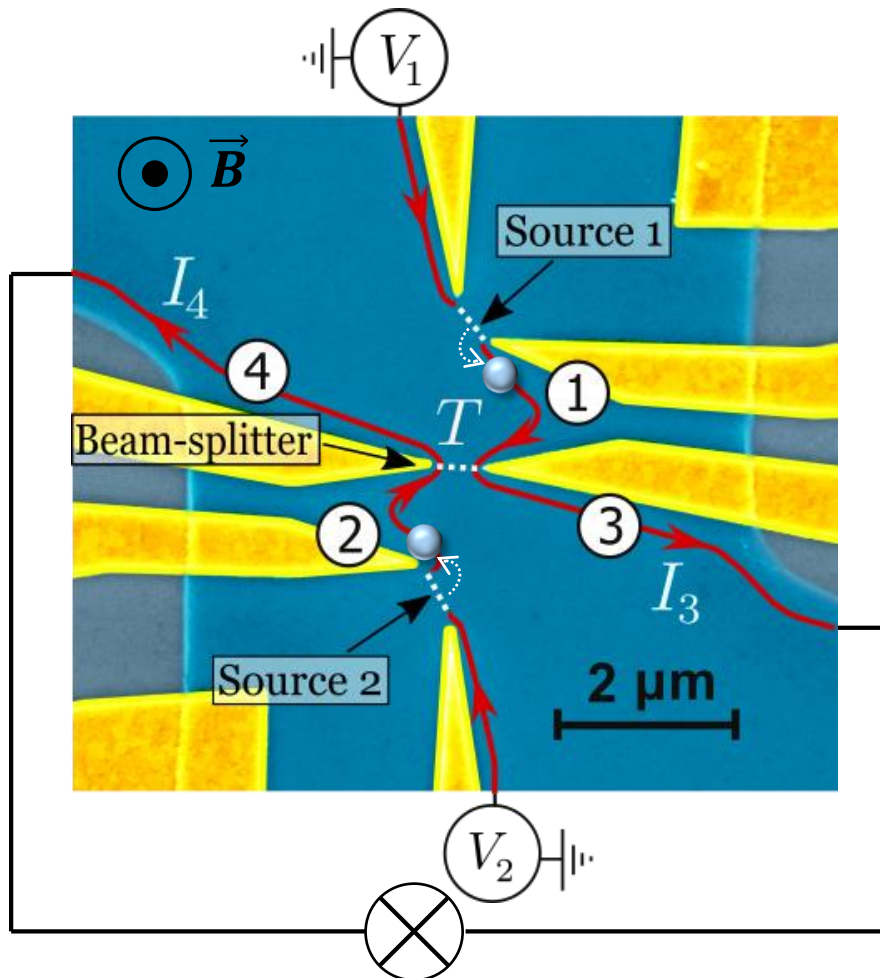
Product of coherences
HBT interferometry



↳ Electronics: $\begin{cases} \psi(t) & \text{Fermion field} \\ I(t) & \text{Electrical current} \end{cases}$

Single and two particle interferometers in quantum Hall conductors

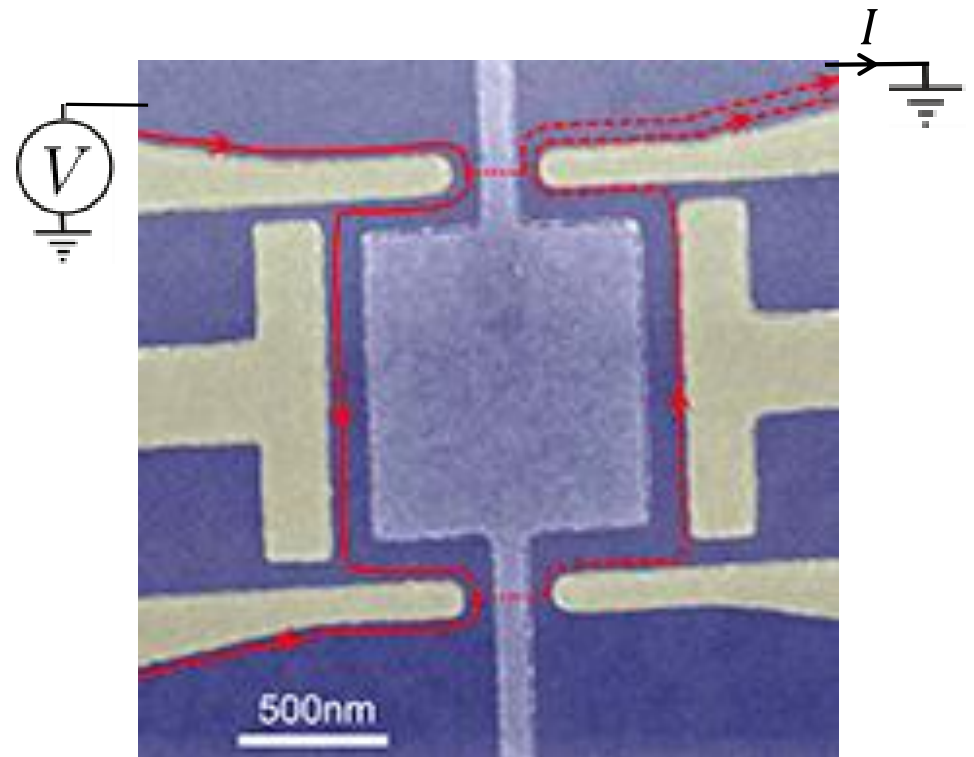
Electron optics experiments in quantum Hall conductors



Current correlations $\langle \delta I_3(t) \delta I_4(t') \rangle$

Two-particle interferometry

H. Bartolomei, M. Kumar et al., Science **368** 173 (2020)



Electrical current $\langle I(t) \rangle$

Single-particle interferometry

J. Nakamura, S. Liang, G.C. Gardner, M.J. Manfra, Nature Physics **16** 931 (2020).

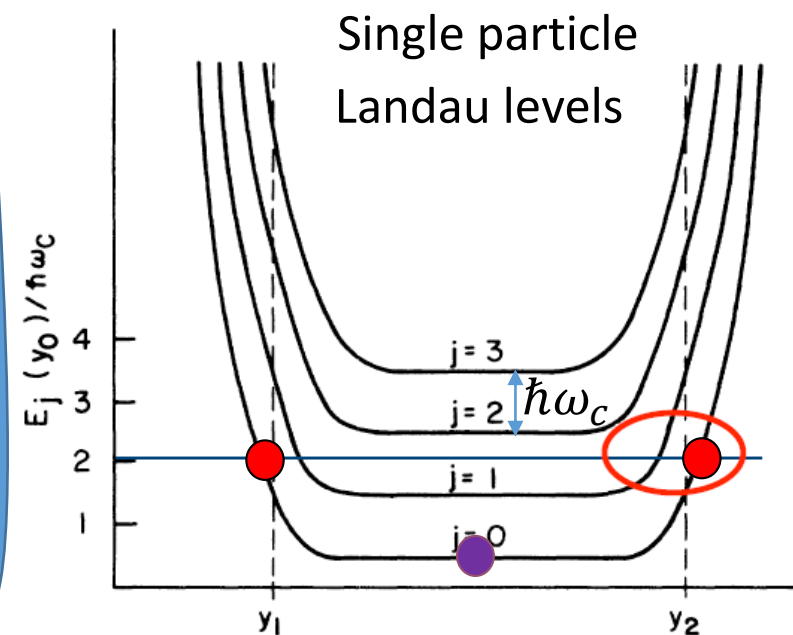
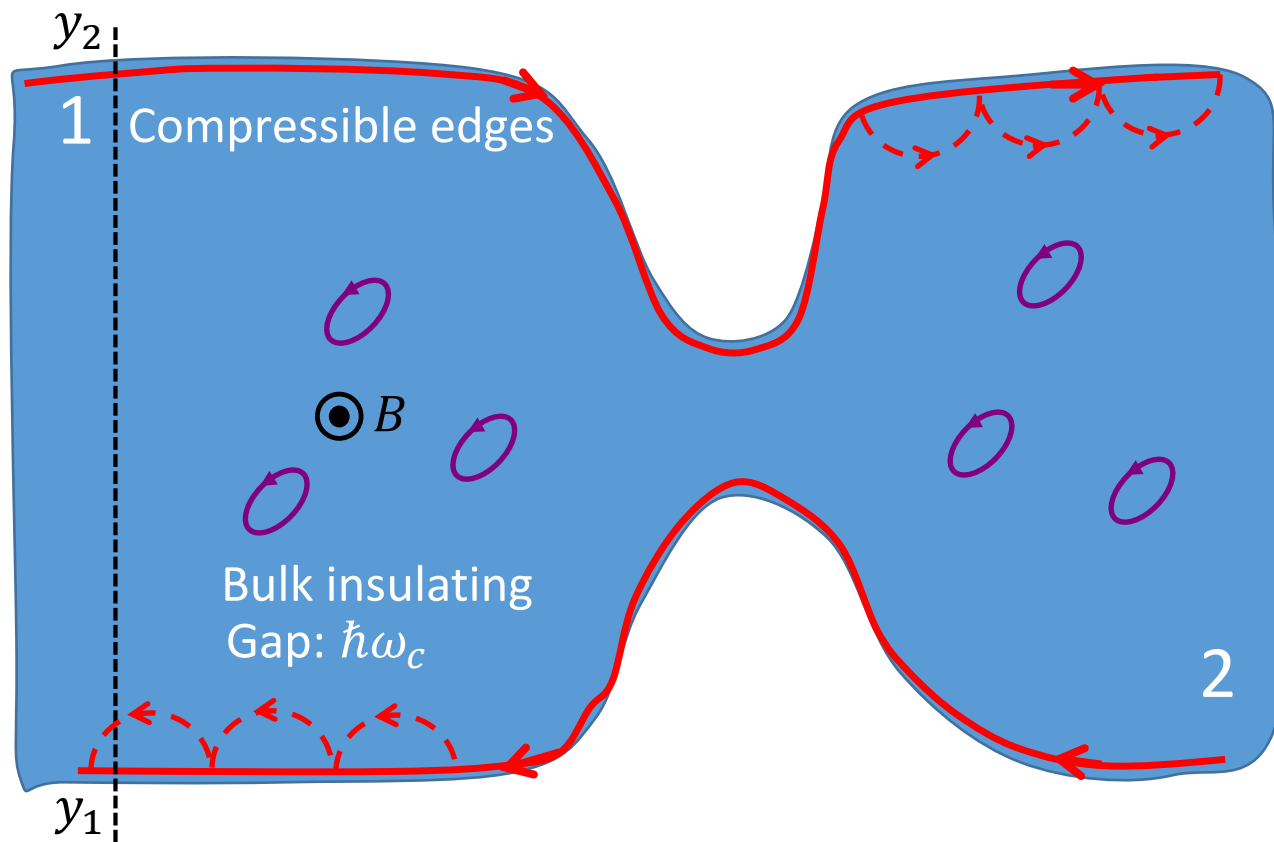
Course 1: Noise and fractional charge

- 1) Electron optics toolbox
- 2) Two particle interferometry, noise and tunneling rates
- 3) Voltage biased quantum point contact: electron and anyon charges

Course 2: noise and fractional statistics

- 1) Electronic/Photonic Hong-Ou-Mandel experiments
- 2) Collider: fermion case
- 3) Collider: anyon case

Integer quantum Hall effect: insulating bulk and transport at the edges



Filling factor: ν (integer) filled Landau levels $\nu = \frac{N_e}{N_s} = \frac{N_e}{\phi/\phi_0} = \frac{N_e}{BSe/h} \searrow$ when $B \nearrow$

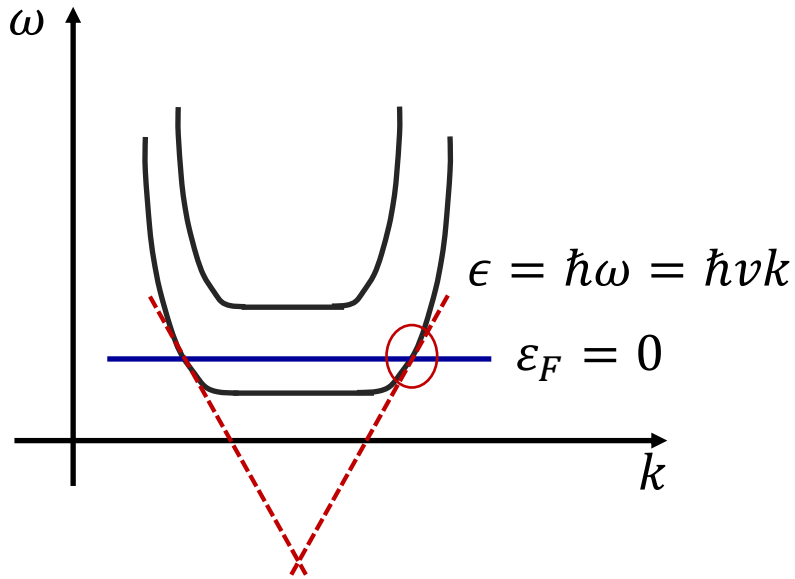
Protection from backscattering

ν 1D wires carrying the current without backscattering $I = \nu \frac{e^2}{h} V$

Electron optics toolbox (Integer case)

Fermion field

Linearization of energy spectrum



- Creation/annihilation operator:

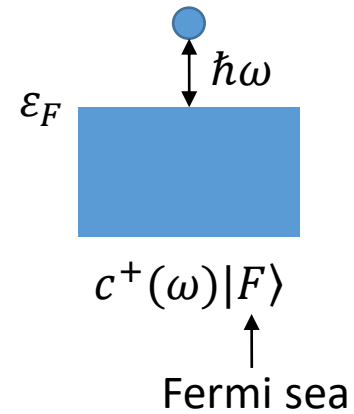
$$\{c(\omega), c^+(\omega')\} = \delta(\omega - \omega')$$

$$H = \int d\omega \hbar\omega c^+(\omega)c(\omega)$$

- Fermion field operator

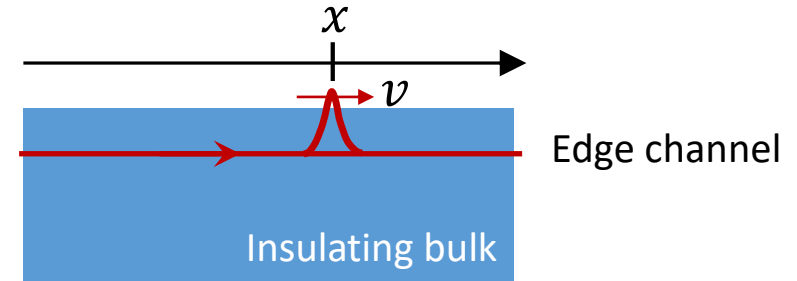
$$\psi(x) = \int \frac{d\omega}{\sqrt{2\pi v}} e^{i\omega x/v} c(\omega)$$

$$\{\psi(x), \psi^+(x')\} = \delta(x - x')$$



- Dynamics of the field: Chiral propagation

$$\psi(x, t) = \int \frac{d\omega}{\sqrt{2\pi v}} e^{i\omega(\frac{x}{v}-t)} c(\omega) = \psi(x - vt)$$



- Current $\frac{\partial \rho}{\partial t} + \frac{\partial I}{\partial x} = 0 \implies I(x, t) = -ev\psi^+(x, t)\psi(x, t)$

In the following
 $v = 1$

- First order coherence/ correlation function of the field

$$G(t, t') = \langle \psi^\dagger(x, t) \psi(x, t') \rangle$$

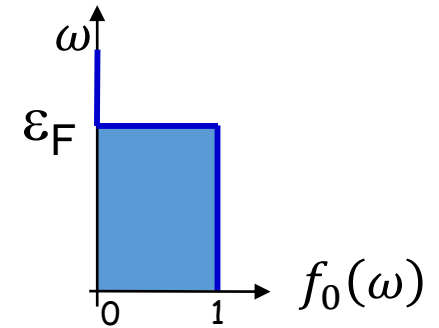
G. Haack et al., Phys. Rev. B **84**, 081303 (2011).

C. Grenier et al., New J. Phys. **13**, 093007 (2011)

Equilibrium correlation function: Fermi sea :

$$G_{eq}(t - t') = \langle F | \psi^\dagger(t) \psi(t') | F \rangle \neq 0$$

$$G_{eq}(t - t') = \int \frac{d\omega}{2\pi} f_0(\omega) e^{i\omega[t-t'-i\tau_c]} = \frac{i/(2\tau_{th})}{\sinh \left[i \frac{\pi\tau_c}{\tau_{th}} - \frac{\pi(t-t')}{\tau_{th}} \right]}$$



$$\tau_{th} = \frac{\hbar}{k_B T_{el}}$$

$$T_{el} = 25mK$$

$$\tau_{th} \approx 300ps$$

$$t - t' \ll \tau_{th}, G_{eq} \approx \frac{1}{2i\pi(t - t' - i\tau_c)}$$

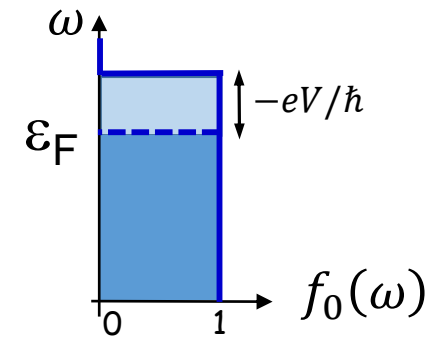
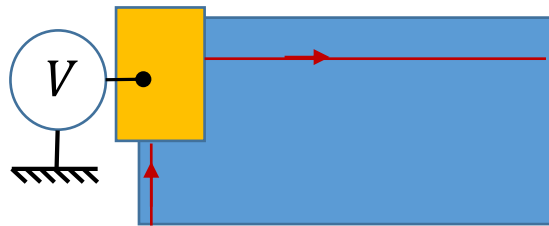
$$t - t' \gg \tau_{th}, G_{eq} \approx e^{-\frac{\pi(t-t')}{\tau_{th}}}$$

- First order coherence/ correlation function of the field

$$G(t, t') = \langle \psi^+(t) \psi(t') \rangle$$

G. Haack et al., Phys. Rev. B **84**, 081303 (2011).

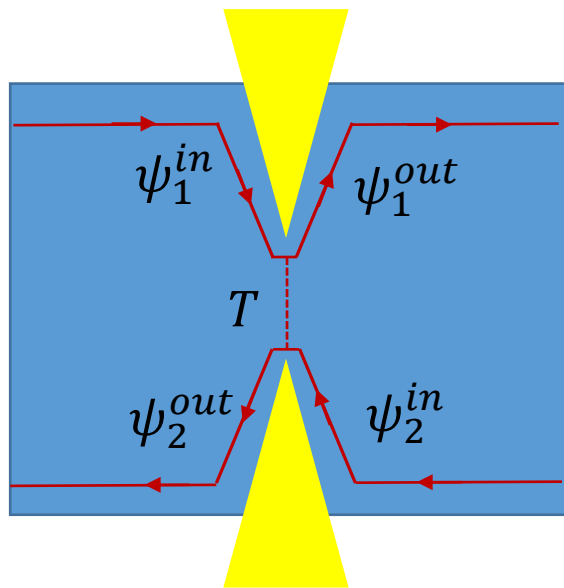
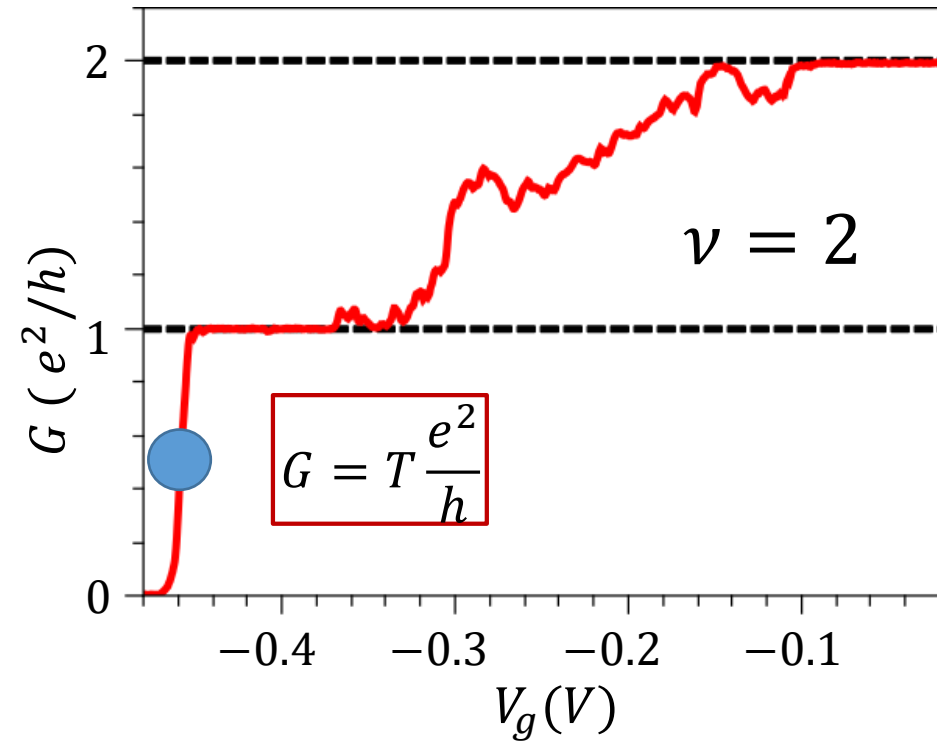
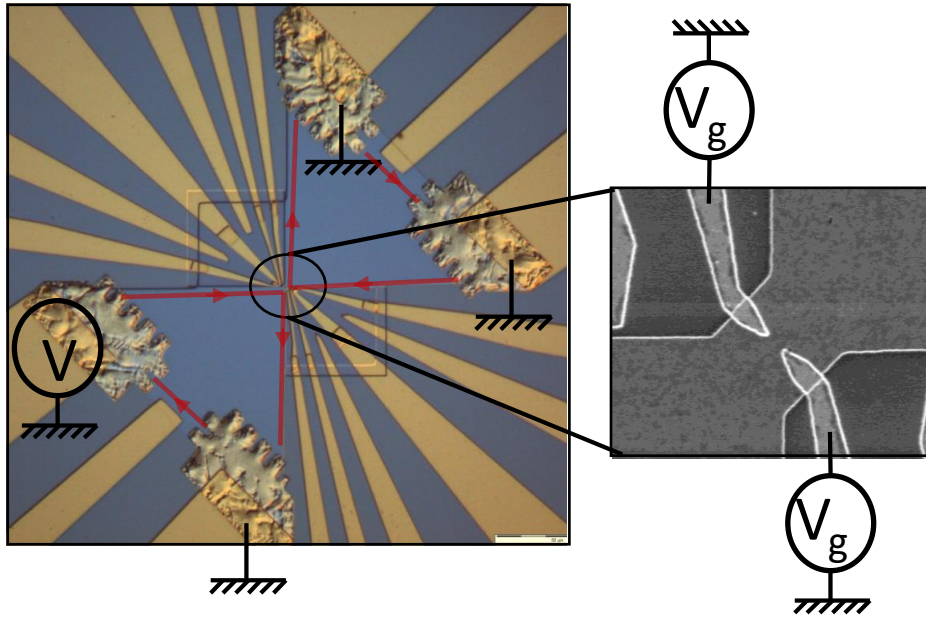
C. Grenier et al., New J. Phys. **13**, 093007 (2011)



$$G_V(t - t') = \int \frac{d\omega}{2\pi} f_0(\omega + eV/\hbar) e^{i\omega[t-t'-i\tau_c]} = e^{-i\frac{eV(t-t')}{\hbar}} G_{eq}(t - t')$$

Electron optics toolbox (IQHE)

The quantum point contact (QPC)

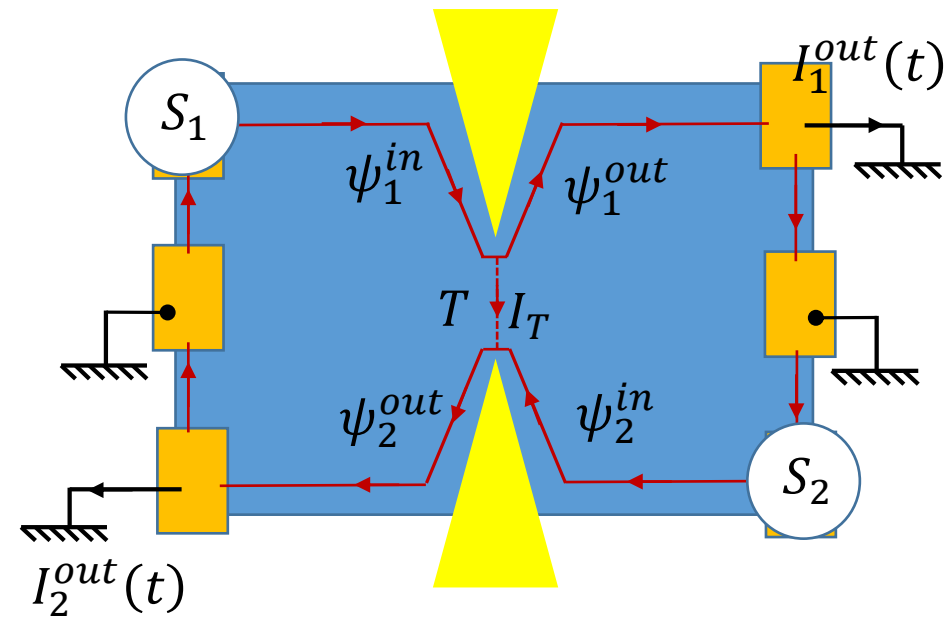
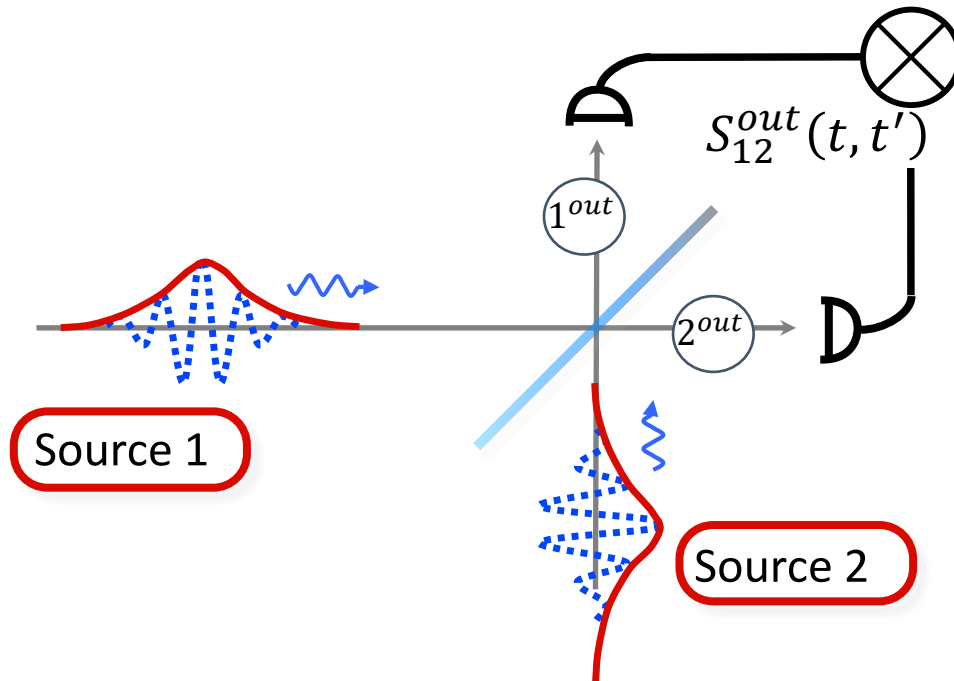


$$\begin{pmatrix} \psi_1^{out} \\ \psi_2^{out} \end{pmatrix} = \underbrace{\begin{pmatrix} r & it \\ it & r \end{pmatrix}}_S \begin{pmatrix} \psi_1^{in} \\ \psi_2^{in} \end{pmatrix}$$

Scattering matrix

$T = |it|^2$: tunneling probability

The electronic Hanbury Brown and Twiss experiment



$$S_{12}^{out}(t, t') = \langle \delta I_1^{out}(t) \delta I_2^{out}(t') \rangle = \langle I_1^{out}(t) I_2^{out}(t') \rangle - \langle I_1^{out}(t) \rangle \langle I_2^{out}(t') \rangle$$

$$I_1^{out}(t) = -e \psi_1^{+,out}(t) \psi_1^{out}(t)$$

$$I_2^{out}(t') = -e \psi_2^{+,out}(t') \psi_2^{out}(t')$$

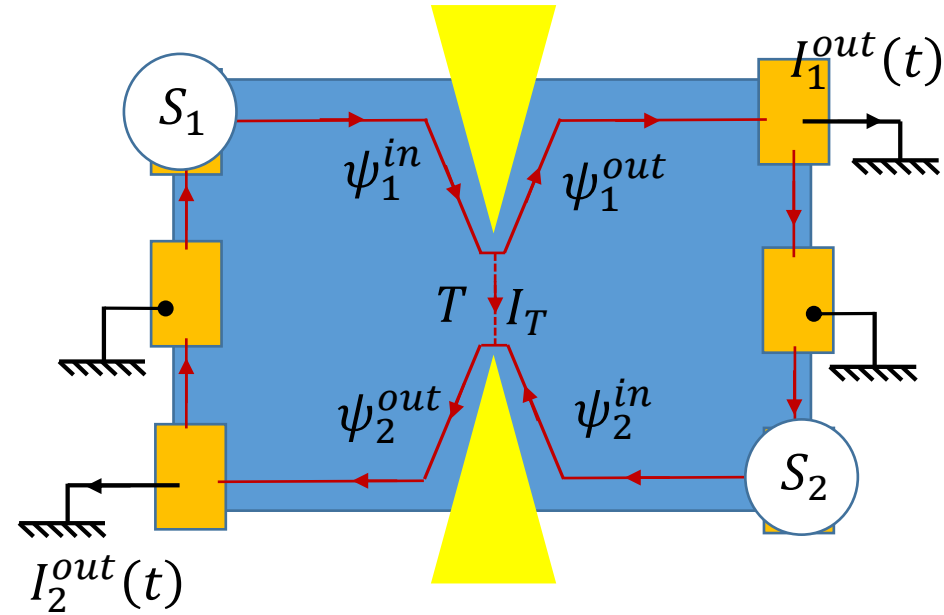
$$\frac{S_{12}^{out}(t, t')}{e^2} = \langle \psi_1^{+,out}(t) \psi_1^{out}(t) \psi_2^{+,out}(t') \psi_2^{out}(t') \rangle - \langle \psi_1^{+,out}(t) \psi_1^{out}(t) \rangle \langle \psi_2^{+,out}(t') \psi_2^{out}(t') \rangle$$

The electronic Hanbury Brown and Twiss experiment

$$\begin{pmatrix} \psi_1^{out} \\ \psi_2^{out} \end{pmatrix} = \begin{pmatrix} r & it \\ it & r \end{pmatrix} \begin{pmatrix} \psi_1^{in} \\ \psi_2^{in} \end{pmatrix}$$

$$R = r^2$$

$$T = t^2$$



$$S_{12}^{out}(t, t') = \langle \delta I_1^{out}(t) \delta I_2^{out}(t') \rangle = \langle I_1^{out}(t) I_2^{out}(t') \rangle - \langle I_1^{out}(t) \rangle \langle I_2^{out}(t') \rangle$$

$$\frac{S_{12}^{out}(t, t')}{e^2} = \langle \psi_1^{+,out}(t) \psi_1^{out}(t) \psi_2^{+,out}(t') \psi_2^{out}(t') \rangle - \langle \psi_1^{+,out}(t) \psi_1^{out}(t) \rangle \langle \psi_2^{+,out}(t') \psi_2^{out}(t') \rangle$$

$$\frac{S_{12}^{out}(t, t')}{e^2} = RT \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle + RT \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle$$

$$+ T^2 \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle + R^2 \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle$$

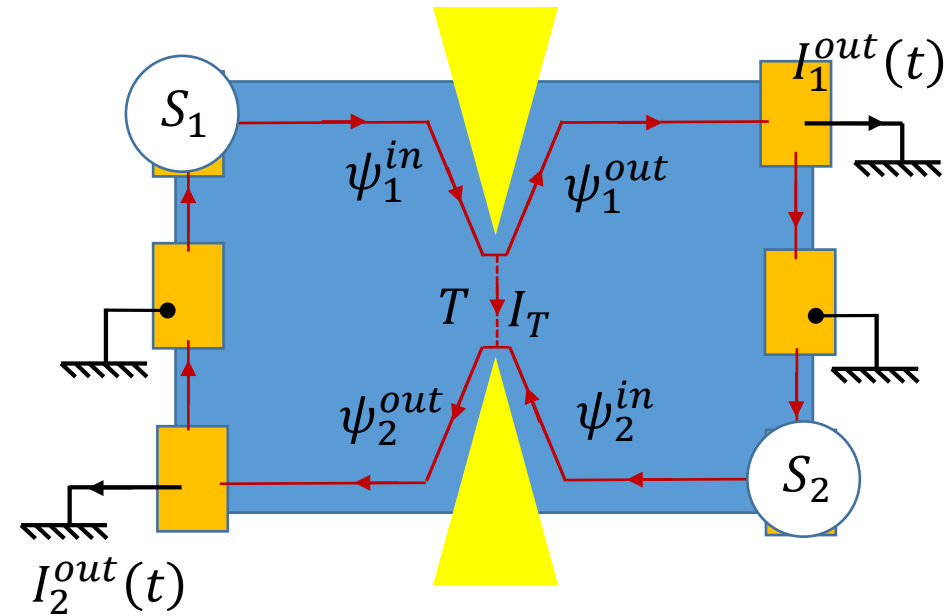
$$- RT \langle \psi_2^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_2^{in}(t') \rangle - RT \langle \psi_1^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_1^{in}(t') \rangle - \sum \langle \quad \rangle \langle \quad \rangle$$

The electronic Hanbury Brown and Twiss experiment

$$\begin{pmatrix} \psi_1^{out} \\ \psi_2^{out} \end{pmatrix} = \begin{pmatrix} r & it \\ it & r \end{pmatrix} \begin{pmatrix} \psi_1^{in} \\ \psi_2^{in} \end{pmatrix}$$

$$R = r^2$$

$$T = t^2$$



$$S_{12}^{out}(t, t') = \langle \delta I_1^{out}(t) \delta I_2^{out}(t') \rangle = \langle I_1^{out}(t) I_2^{out}(t') \rangle - \langle I_1^{out}(t) \rangle \langle I_2^{out}(t') \rangle$$

$$\frac{S_{12}^{out}(t, t')}{e^2} = \langle \psi_1^{+,out}(t) \psi_1^{out}(t) \psi_2^{+,out}(t') \psi_2^{out}(t') \rangle - \langle \psi_1^{+,out}(t) \psi_1^{out}(t) \rangle \langle \psi_2^{+,out}(t') \psi_2^{out}(t') \rangle$$

$$\frac{S_{12}^{out}(t, t')}{e^2} = RT \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle + RT \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle$$

$$+ T^2 \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle + R^2 \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle$$

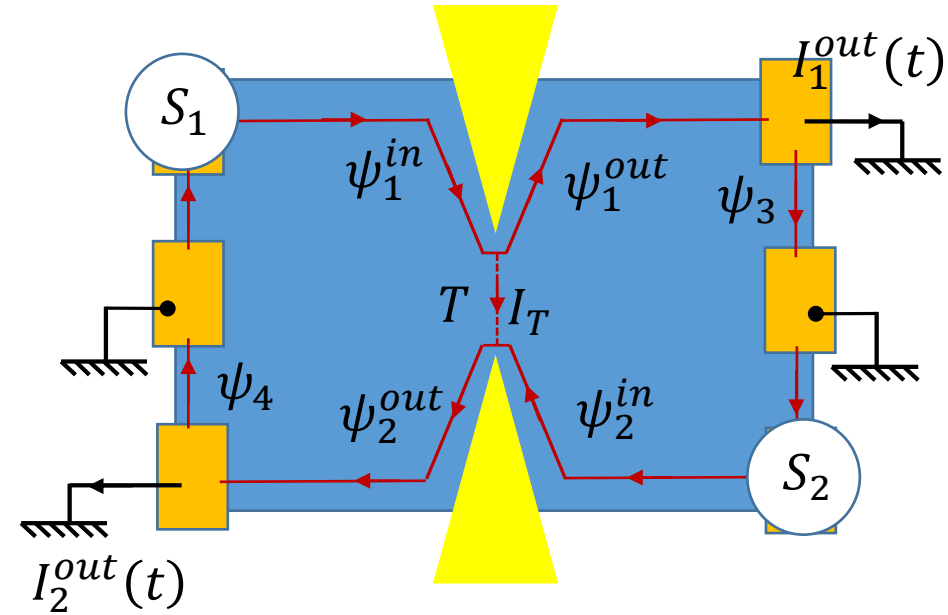
$$- RT \langle \psi_2^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_2^{in}(t') \rangle - RT \langle \psi_1^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_1^{in}(t') \rangle - \sum \langle \rangle \langle \rangle$$

The electronic Hanbury Brown and Twiss experiment

$$\begin{pmatrix} \psi_1^{out} \\ \psi_2^{out} \end{pmatrix} = \begin{pmatrix} r & it \\ it & r \end{pmatrix} \begin{pmatrix} \psi_1^{in} \\ \psi_2^{in} \end{pmatrix}$$

$$R = r^2$$

$$T = t^2$$



$$S_{12}^{out}(t, t') = \langle \delta I_1^{out}(t) \delta I_2^{out}(t') \rangle = \langle I_1^{out}(t) I_2^{out}(t') \rangle - \langle I_1^{out}(t) \rangle \langle I_2^{out}(t') \rangle$$

$$\frac{S_{12}^{out}(t, t')}{e^2} = RT \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle - RT \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \rangle \langle \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle$$

$$+ RT \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle - RT \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \rangle \langle \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle$$

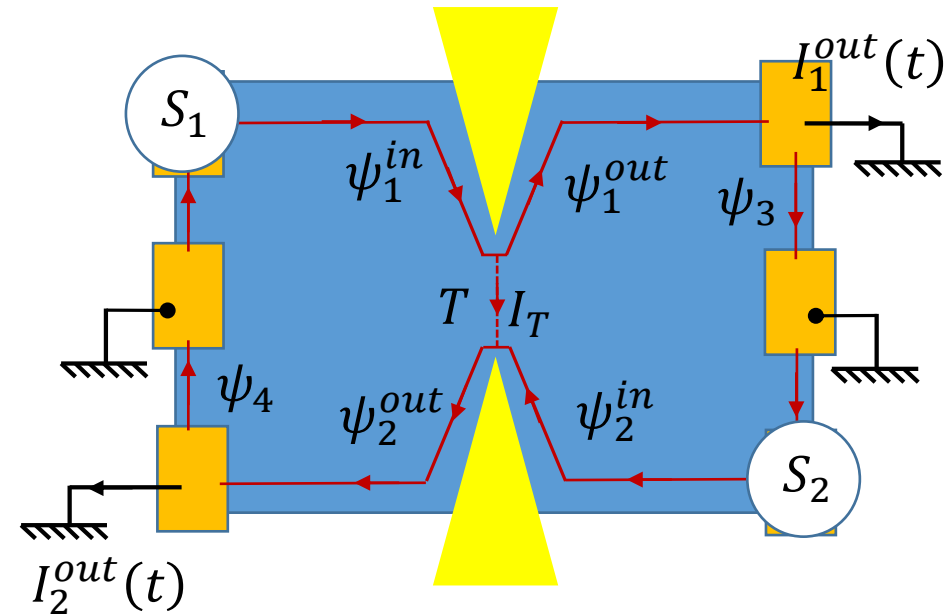
$$- RT \left[\langle \psi_2^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_2^{in}(t') \rangle + \langle \psi_1^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_1^{in}(t') \rangle \right]$$

The electronic Hanbury Brown and Twiss experiment

$$\begin{pmatrix} \psi_1^{out} \\ \psi_2^{out} \end{pmatrix} = \begin{pmatrix} r & it \\ it & r \end{pmatrix} \begin{pmatrix} \psi_1^{in} \\ \psi_2^{in} \end{pmatrix}$$

$$R = r^2$$

$$T = t^2$$



$$S_{12}^{out}(t, t') = \langle \delta I_1^{out}(t) \delta I_2^{out}(t') \rangle = \langle I_1^{out}(t) I_2^{out}(t') \rangle - \langle I_1^{out}(t) \rangle \langle I_2^{out}(t') \rangle$$

$$RTS_{11}^{in}$$

$$\frac{S_{12}^{out}(t, t')}{e^2} = \cancel{RT \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle - RT \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \rangle \langle \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle}$$

$$+ \cancel{RT \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle - RT \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \rangle \langle \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle} \quad RTS_{22}^{in}$$

$$- RT \left[\langle \psi_2^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_2^{in}(t') \rangle + \langle \psi_1^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_1^{in}(t') \rangle \right]$$

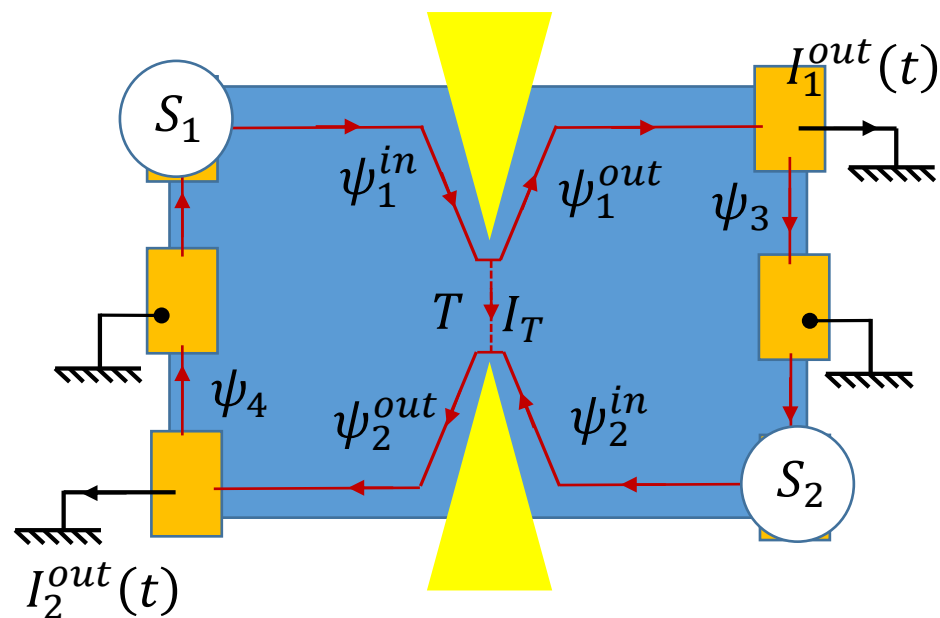
$$-S_{I_T}$$

The electronic Hanbury Brown and Twiss experiment, weak backscattering $T \ll 1$

$$\begin{pmatrix} \psi_1^{out} \\ \psi_2^{out} \end{pmatrix} = \begin{pmatrix} r & it \\ it & r \end{pmatrix} \begin{pmatrix} \psi_1^{in} \\ \psi_2^{in} \end{pmatrix}$$

$$R = r^2$$

$$T = t^2$$



$$S_{12}^{out}(t, t') = \langle \delta I_1^{out}(t) \delta I_2^{out}(t') \rangle = \langle I_1^{out}(t) I_2^{out}(t') \rangle - \langle I_1^{out}(t) \rangle \langle I_2^{out}(t') \rangle$$

$$RTS_{11}^{in} \approx TS_{11}^{in} \quad (T \ll 1)$$

$$\frac{S_{12}^{out}(t, t')}{e^2} = T \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle - RT \langle \psi_1^{+,in}(t) \psi_1^{in}(t) \rangle \langle \psi_1^{+,in}(t') \psi_1^{in}(t') \rangle$$

$$+ T \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle - RT \langle \psi_2^{+,in}(t) \psi_2^{in}(t) \rangle \langle \psi_2^{+,in}(t') \psi_2^{in}(t') \rangle \approx TS_{22}^{in}$$

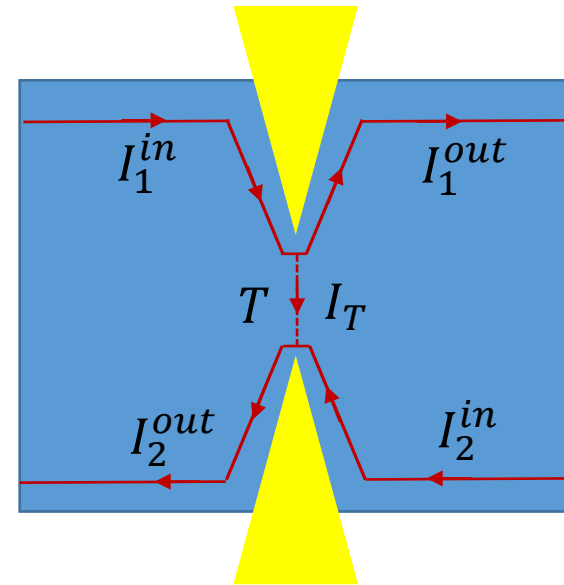
$$- T \left[\langle \psi_2^{+,in}(t) \psi_1^{in}(t) \psi_1^{+,in}(t') \psi_2^{in}(t') \rangle + \langle \psi_1^{+,in}(t) \psi_2^{in}(t) \psi_2^{+,in}(t') \psi_1^{in}(t') \rangle \right]$$

$$-S_{I_T}$$

Cross-correlations S_{12}^{out} and noise of the tunneling current S_{I_T}

$$I_1^{out} = I_1^{in} - I_T$$

$$I_2^{out} = I_2^{in} + I_T$$

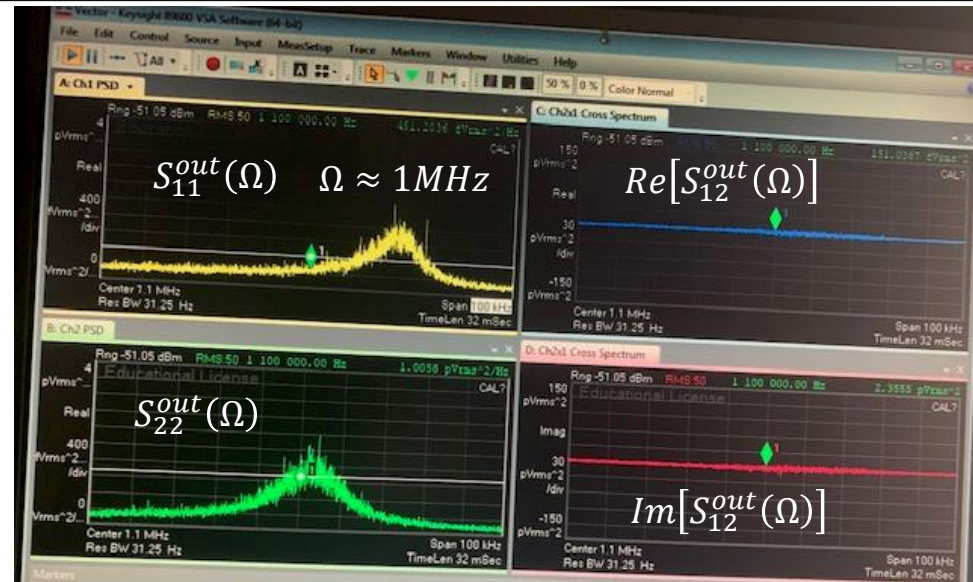
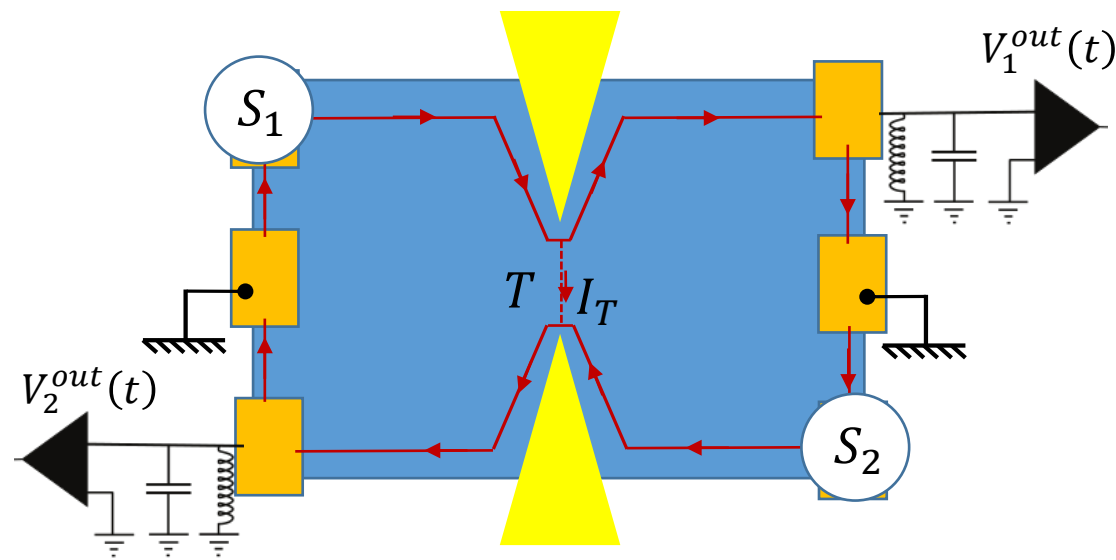


$$\langle \delta I_1^{out} \delta I_2^{out} \rangle = \langle \delta I_1^{in} \delta I_2^{in} \rangle + \langle \delta I_1^{in} \delta I_T \rangle - \langle \delta I_T \delta I_2^{in} \rangle - \langle \delta I_T \delta I_T \rangle$$

$$\delta I_T \approx T[\delta I_1^{in} - \delta I_2^{in}]$$

$$\langle \delta I_1^{out} \delta I_2^{out} \rangle = T \langle \delta I_1^{in} \delta I_1^{in} \rangle + T \langle \delta I_2^{in} \delta I_2^{in} \rangle - \langle \delta I_T \delta I_T \rangle$$

$$S_{12}^{out} = T[S_{11}^{in} + S_{22}^{in}] - S_{I_T}$$



$$S_{12}^{out}(\Omega) = 2 \lim_{T_{meas} \rightarrow \infty} \frac{1}{T_{meas}} \int d\bar{t} d\tau S_{12}^{out}\left(\bar{t} + \frac{\tau}{2}, \bar{t} - \frac{\tau}{2}\right) e^{i\Omega\tau} = \frac{2 \langle \delta I_1^{out}(\Omega) \delta I_2^{out}(-\Omega) \rangle}{T_{meas}}$$

$$\text{Low frequency noise: } S_{12}^{out}(\Omega = 0) = 2 \lim_{T_{meas} \rightarrow \infty} \frac{1}{T_{meas}} \int d\bar{t} d\tau S_{12}^{out}\left(\bar{t} + \frac{\tau}{2}, \bar{t} - \frac{\tau}{2}\right)$$

Stationary case: no dependence on \bar{t} :

$$S_{I_T}(\Omega = 0) = 2e^2 \left[\underbrace{T \int d\tau \langle \psi_1^{+,in}(\tau) \psi_1^{in}(0) \rangle \langle \psi_2^{in}(\tau) \psi_2^{+,in}(0) \rangle}_{\Gamma_{1 \rightarrow 2}} + T \int d\tau \langle \psi_1^{in}(\tau) \psi_1^{+,in}(0) \rangle \langle \psi_2^{+,in}(\tau) \psi_2^{in}(0) \rangle \right]_{\Gamma_{2 \rightarrow 1}}$$

Example: voltage biased QPC

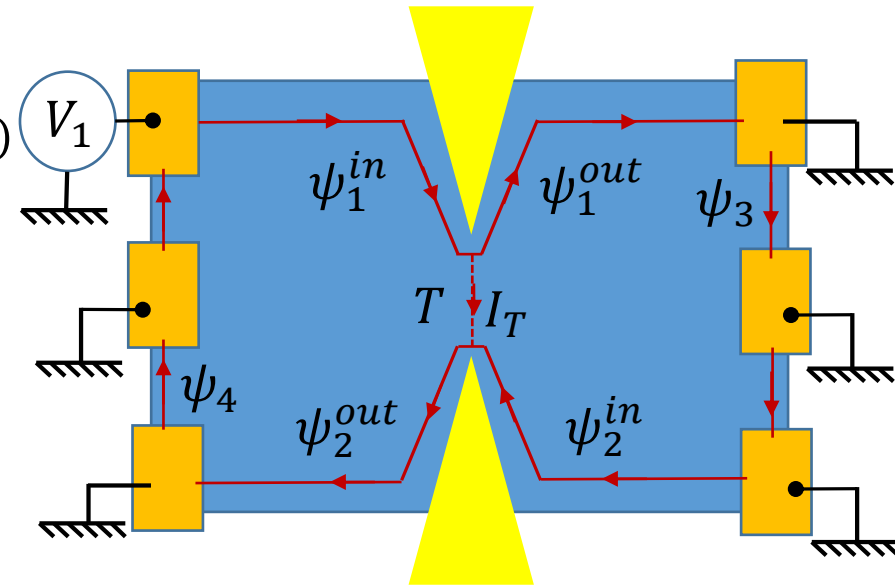
$$\langle \psi_1^{+,in}(\tau) \psi_1^{in}(0) \rangle = e^{-i\frac{eV_1\tau}{\hbar}} G_{eq}(\tau) = \int \frac{d\omega}{2\pi} e^{i(\omega - eV_1/\hbar)\tau} f_0(\omega)$$

$$\langle \psi_1^{in}(\tau) \psi_1^{+,in}(0) \rangle = e^{i\frac{eV_1\tau}{\hbar}} G_{eq}(\tau) = \int \frac{d\omega}{2\pi} e^{i(\omega + eV_1/\hbar)\tau} f_0(\omega)$$

$$\Gamma_{1 \rightarrow 2} = T \int \frac{d\omega d\omega'}{2\pi} \delta(\omega + \omega' - eV_1/\hbar) f_0(\omega) f_0(\omega')$$

$$\Gamma_{1 \rightarrow 2} = T \int \frac{d\omega}{2\pi} f_0(\omega + eV_1/\hbar) [1 - f_0(\omega)]$$

$$\Gamma_{2 \rightarrow 1} = T \int \frac{d\omega}{2\pi} [1 - f_0(\omega + eV_1/\hbar)] f_0(\omega)$$

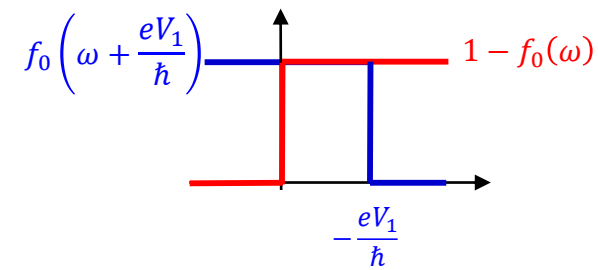


Classical picture: $\frac{eV_1\tau_{th}}{\hbar} \gg 1$ (zero temperature limit)

$$V_1 < 0 \quad \Gamma_{1 \rightarrow 2} = -T \frac{eV_1}{\hbar} \quad \Gamma_{2 \rightarrow 1} = 0$$

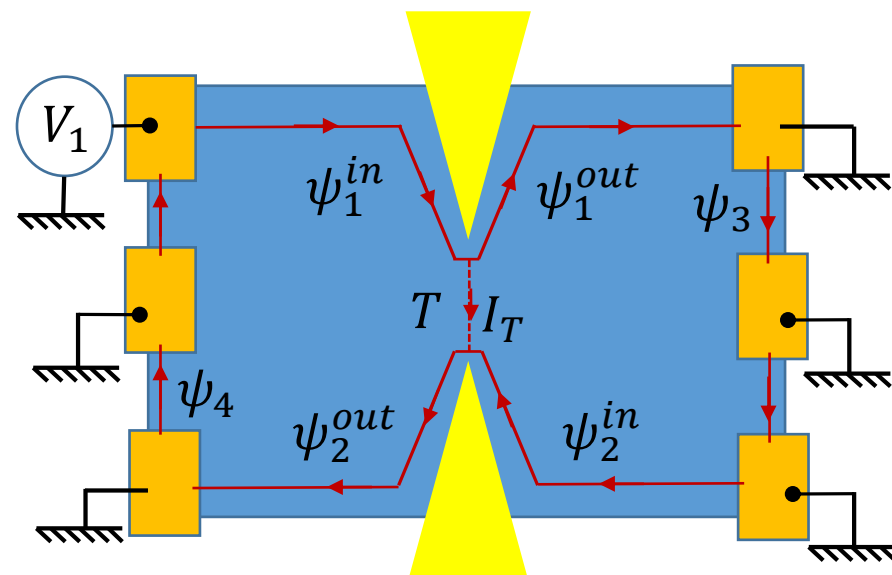
$$\Gamma_{1 \rightarrow 2} = T\Gamma_1^{in} \quad \text{Tunelling rate for electrons from } 1 \rightarrow 2$$

$$V_1 > 0 \quad \Gamma_{1 \rightarrow 2} = 0 \quad \Gamma_{2 \rightarrow 1} = T\Gamma_2^{in} \quad \text{Tunelling rate for } e^- \text{ from } 2 \rightarrow 1$$



$$\Gamma_{1 \rightarrow 2} = T \int \frac{d\omega}{2\pi} f_0(\omega + eV_1/\hbar) [1 - f_0(\omega)]$$

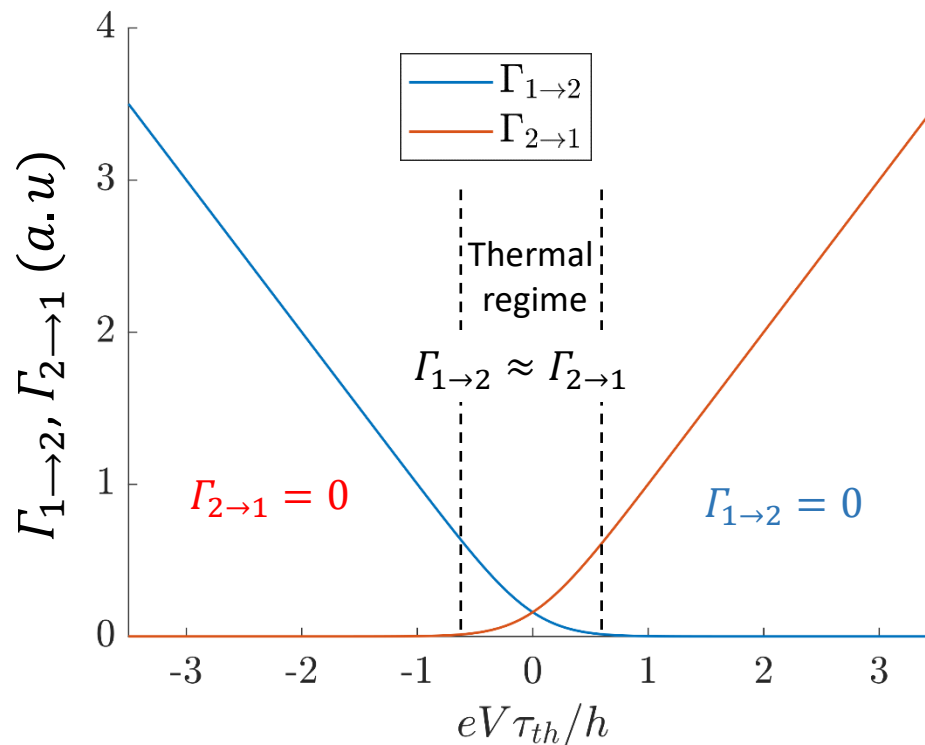
$$\Gamma_{2 \rightarrow 1} = T \int \frac{d\omega}{2\pi} [1 - f_0(\omega + eV_1/\hbar)] f_0(\omega)$$



$$z = \frac{eV\tau_{th}}{h}$$

$$\Gamma_{1 \rightarrow 2} = -\frac{T}{\tau_{th}} z \frac{e^{-2\pi z}}{e^{-2\pi z} - 1}$$

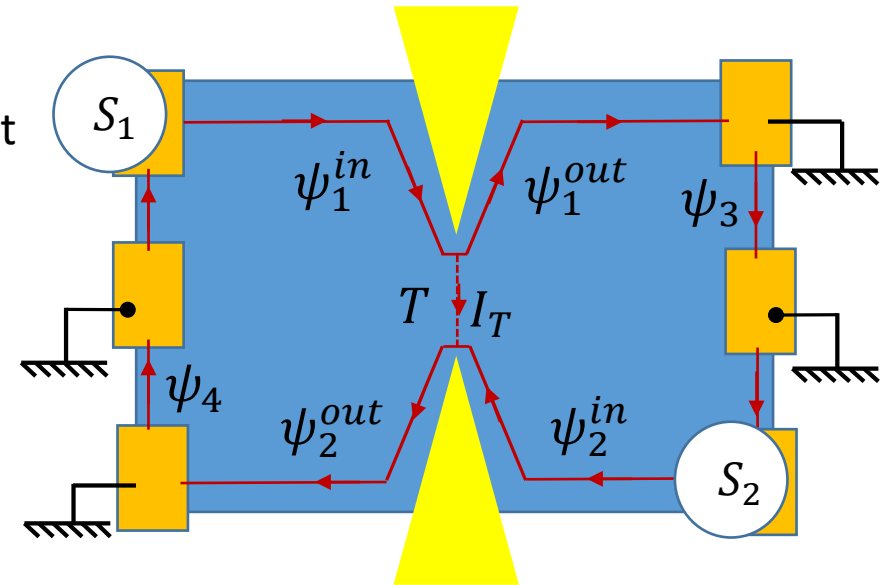
$$\Gamma_{2 \rightarrow 1} = \frac{T}{\tau_{th}} z \frac{e^{2\pi z}}{e^{2\pi z} - 1}$$



Result for out of equilibrium tunneling at the beam-splitter at first order in the tunneling Hamiltonian H_T :

$$H_T = \zeta \psi_1^+ \psi_2 + \zeta^* \psi_2^+ \psi_1$$

Valid in the integer and fractional quantum Hall regime



Observable quantities:

$$\langle I_T \rangle = e(\Gamma_{2 \rightarrow 1} - \Gamma_{1 \rightarrow 2})$$

$$S_{12}^{out} = T[S_{11}^{in} + S_{22}^{in}] - S_{I_T}$$

$$S_{I_T} = 2e^2(\Gamma_{2 \rightarrow 1} + \Gamma_{1 \rightarrow 2})$$

$$\Gamma_{1 \rightarrow 2} = T \int d\tau \langle \psi_1^{+,in}(\tau) \psi_1^{in}(0) \rangle \langle \psi_2^{in}(\tau) \psi_2^{+,in}(0) \rangle$$

$$\Gamma_{2 \rightarrow 1} = T \int d\tau \langle \psi_2^{+,in}(\tau) \psi_2^{in}(0) \rangle \langle \psi_1^{in}(\tau) \psi_1^{+,in}(0) \rangle$$

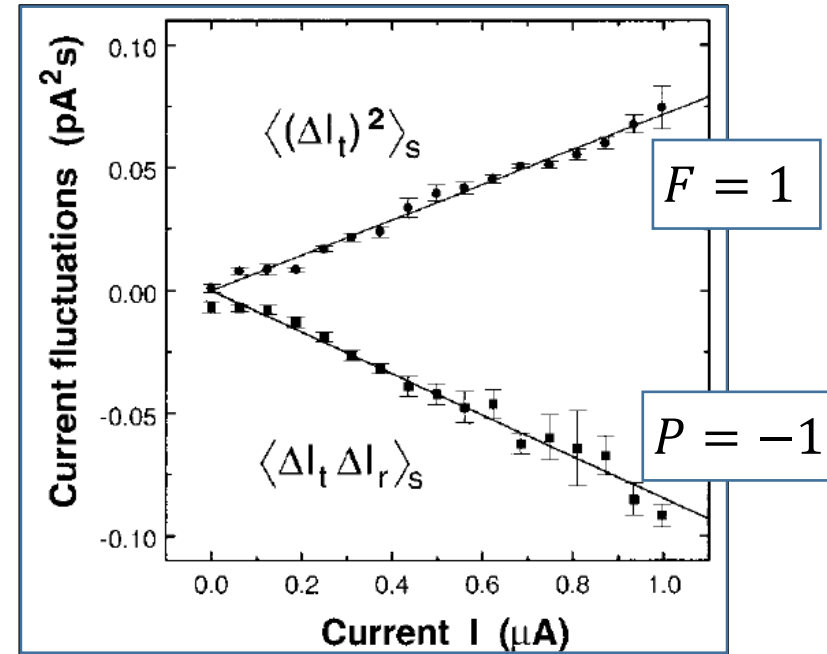
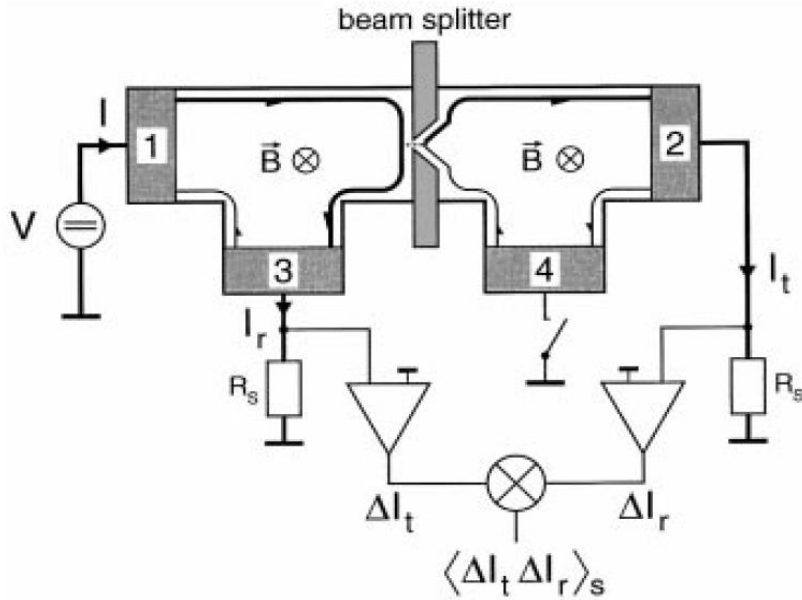
Fano factors:

$$F = \frac{S_{I_T}}{2eI_T} = \frac{\Gamma_{2 \rightarrow 1} + \Gamma_{1 \rightarrow 2}}{\Gamma_{2 \rightarrow 1} - \Gamma_{1 \rightarrow 2}}$$

DC biased QPC: $F = 1$

$$P = \frac{S_{12}^{out}}{2eT(I_1 + I_2)} = \frac{S_{12}^{out}}{2eTI_+}$$

DC biased QPC: $P = -1$



M. Henny et al., Science **284** 296 (1999)

$$S_{11}^{in} = S_{22}^{in} = 0 \quad \rightarrow \quad S_{11}^{out} = S_{22}^{out} = -S_{12}^{out} = S_{I_T}$$

Single source dc case:

$$T \ll 1$$

Weak-backscattering

$$F = \frac{S_{I_T}}{2e^* I_T}$$

$$P = \frac{S_{12}^{out}}{2e T I_1}$$

$$I_T = T I_1$$

$$P = -F = -1$$

Orders of magnitude:

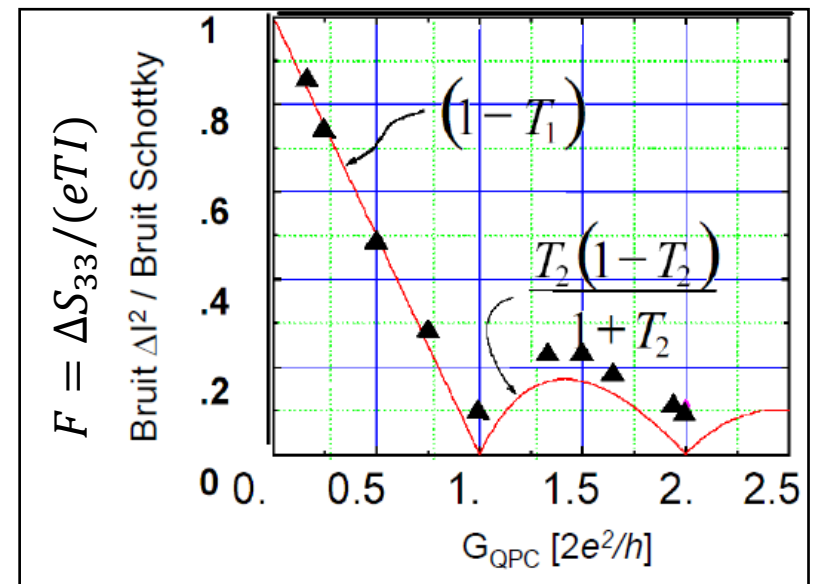
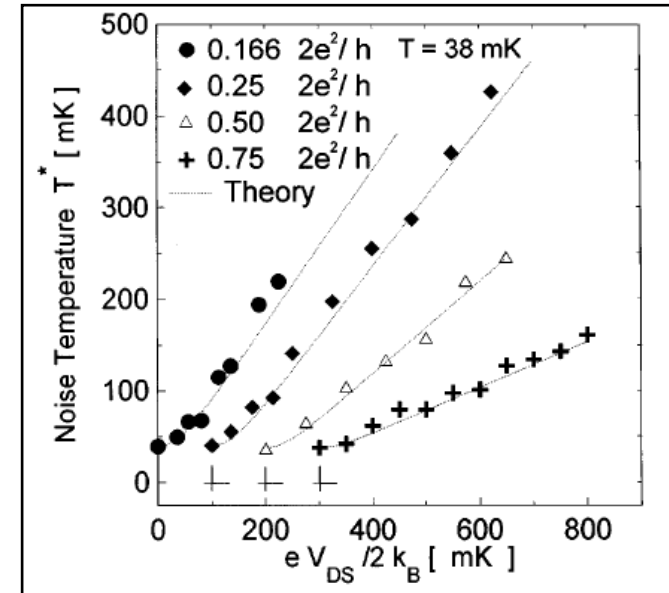
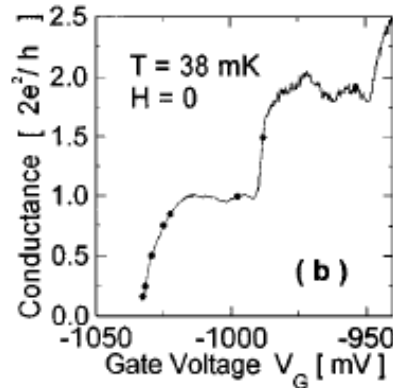
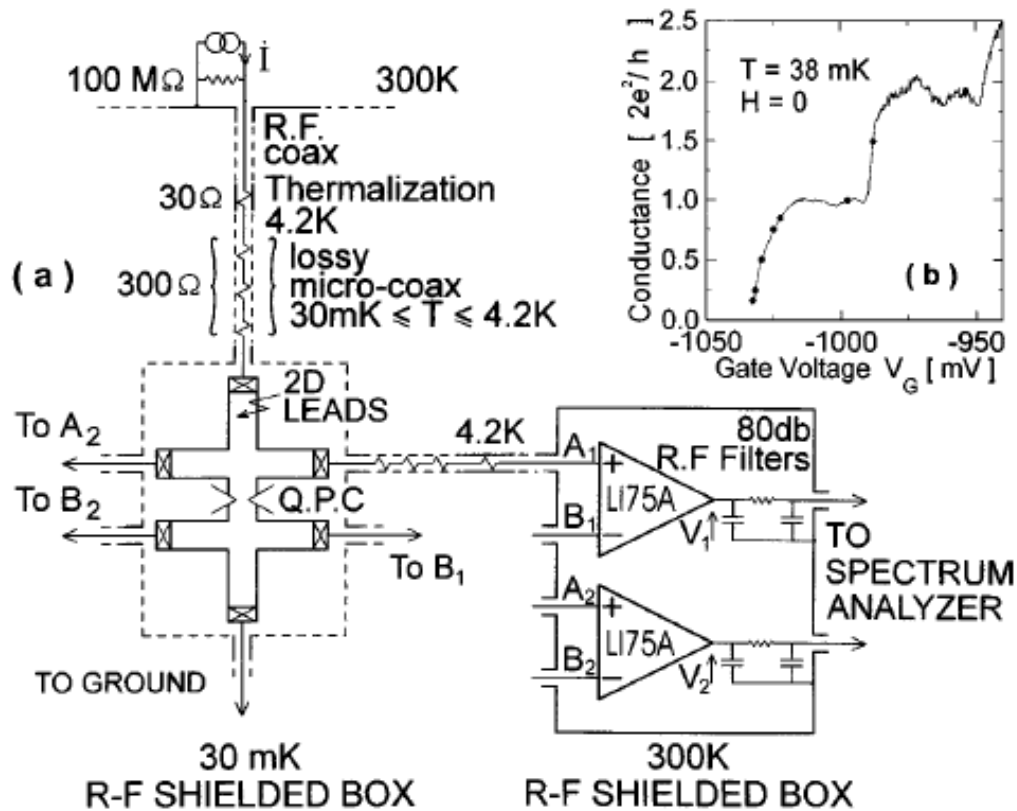
$$V \leq 100 \mu V \quad \frac{eV}{k_B} \leq 1 K$$

$$S_I \leq 2eT \frac{e^2}{h} V \approx qq 10^{-28} A^2 \cdot Hz^{-1} \quad (T = 0.1) \rightarrow T_N \approx 20 mK$$

Cold amplifiers: $T_{N,a} \approx 100 mK$

100 μK resolution in \sim minute

First measurements of partition noise in quantum point contact



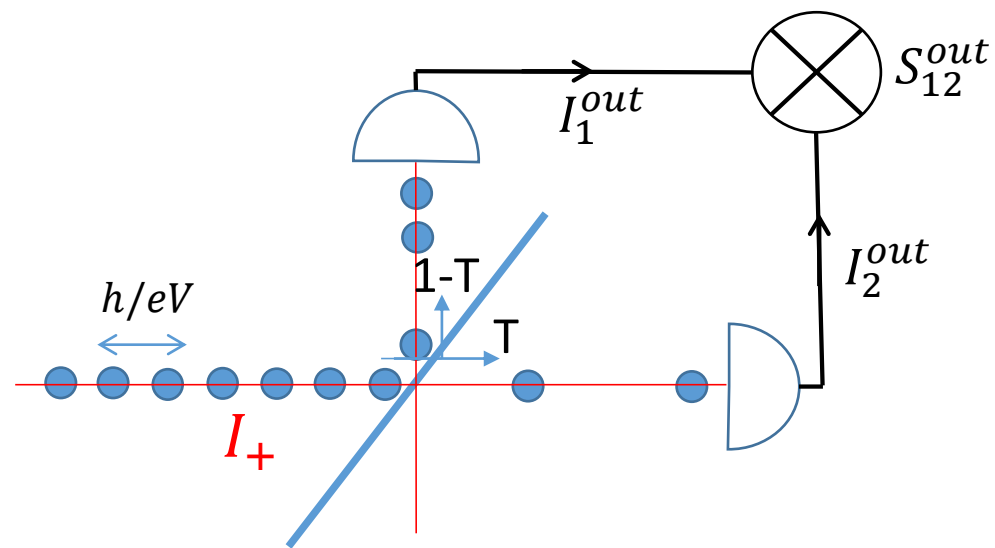
Deviations from weak-backscattering regime:

$$S_{IT} \propto RT \dots = T(1 - T) \dots$$

$$F = -P = (1 - T)$$

A. Kumar et al.,
Phys. Rev. Lett. **76** 2778 (1996)

Single quantum point contact in the dc regime: electron charge

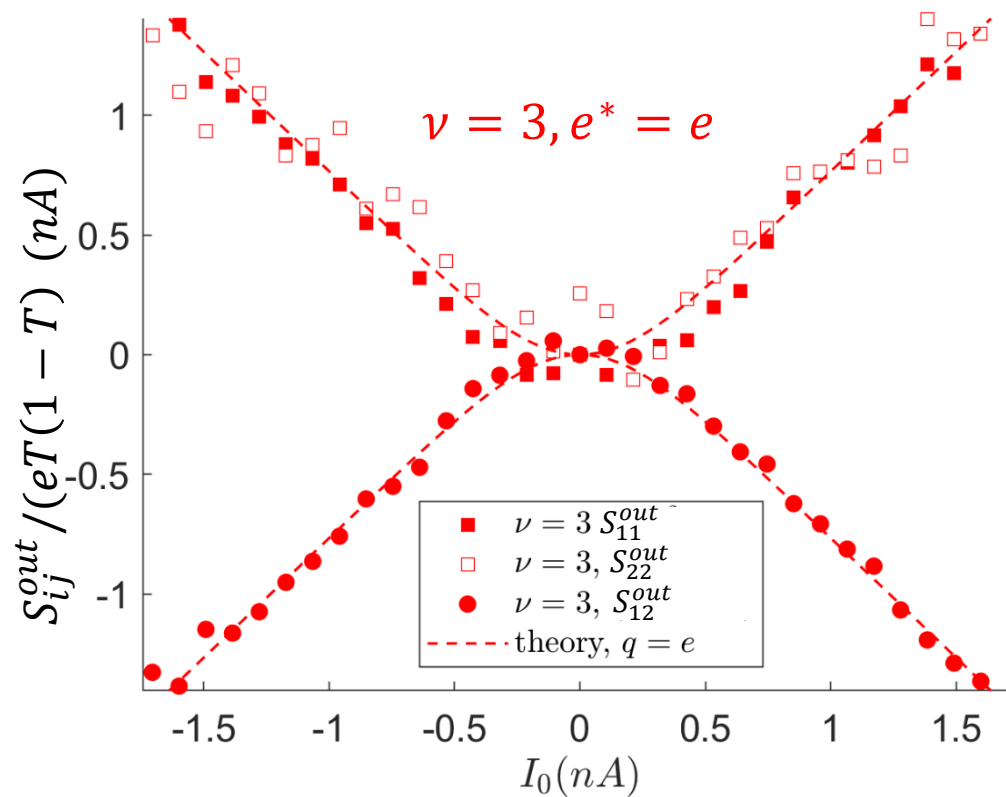
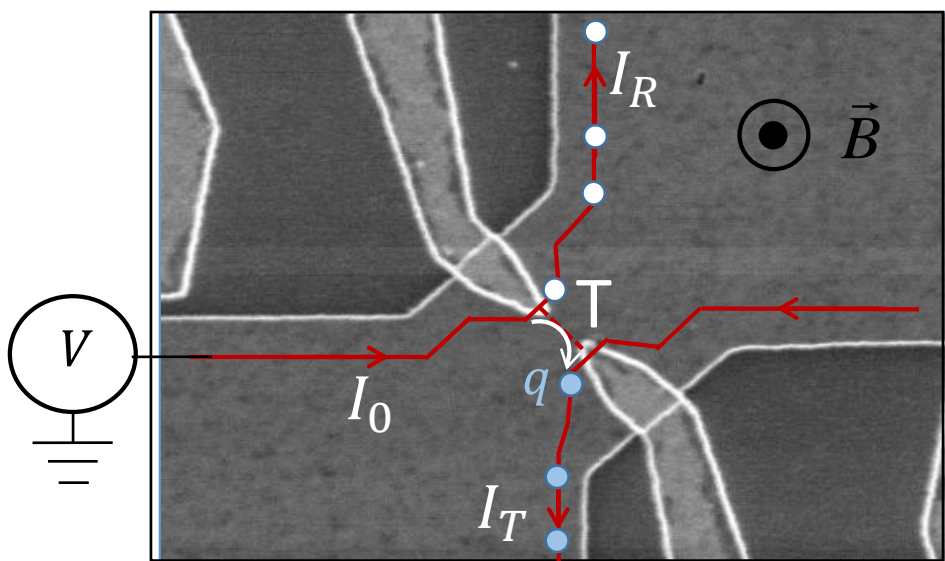


Binomial law: $\langle \Delta N_T^2 \rangle = T(1 - T) N_+$

$$S_{rp} = 2eT(1 - T)I_+$$

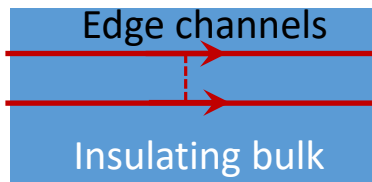
$$S_{rp} \approx 2eTI_+ \quad (T \ll 1)$$

$$S_{12}^{out} = -S_{11}^{out} = -S_{22}^{out} = S_{rp}$$



$$T_{el} \approx 30 \text{ mK}$$

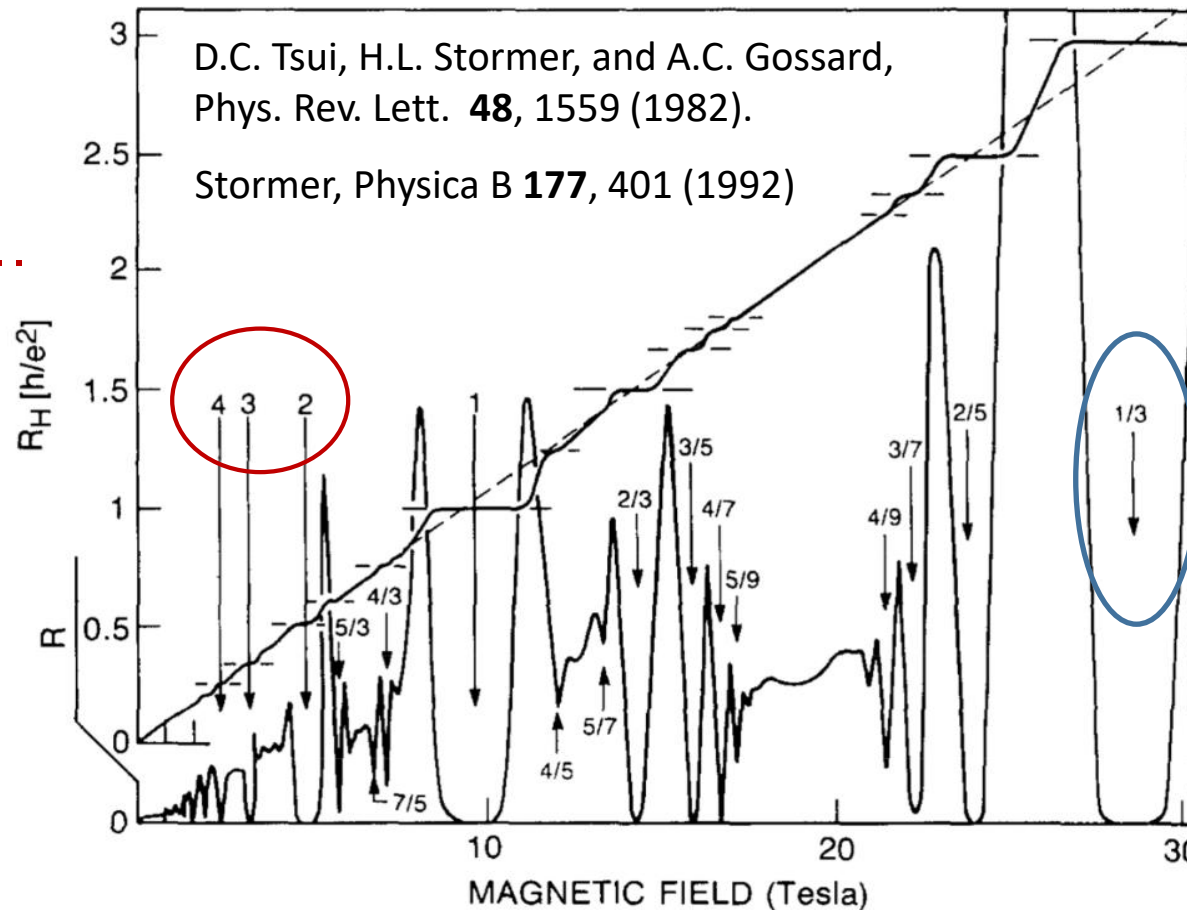
IQHE, $\nu = 1, 2, 3, \dots$



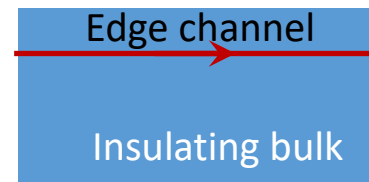
$$G_0 = N \frac{e^2}{h}$$

$$e^* = e$$

$$\varphi = \pi$$



FQHE, $\nu = 1/3$



$$G_0 = \frac{1}{3} \frac{e^2}{h}$$

$$e^* = e/3$$

$$\varphi = \pi/3$$

Each FQHE phase hosts a specific variety of anyons characterized by their fractional charge q and their fractional statistics φ

Halperin, PRL 52 1583 (1984)

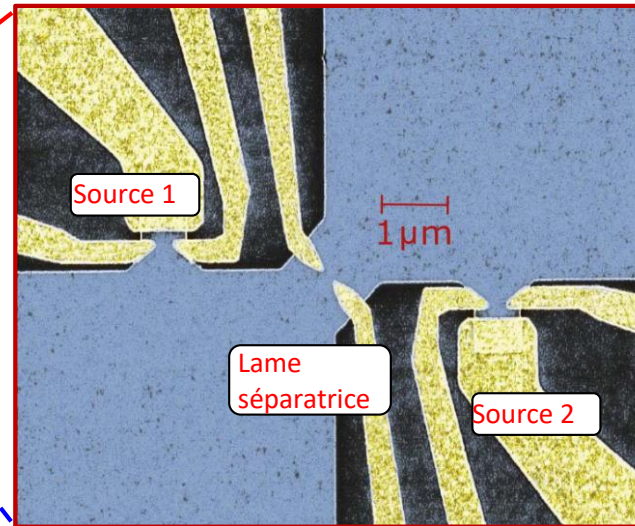
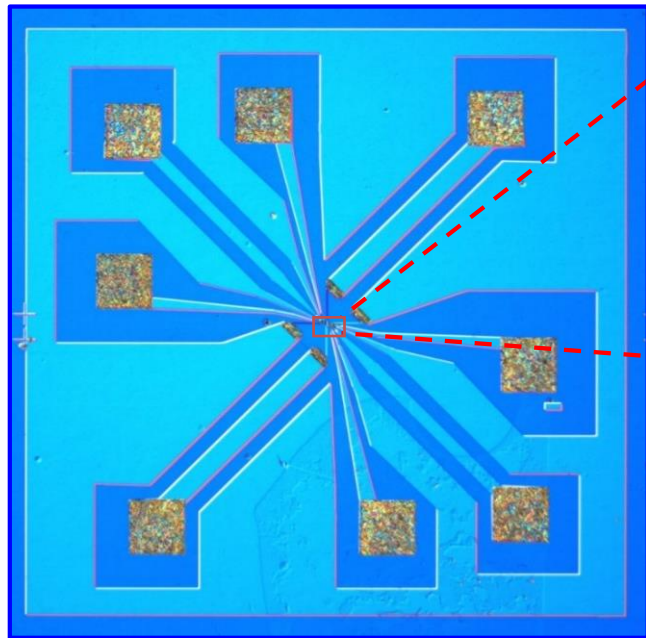
Arovas, Schrieffer, Wilczek PRL 53 722 (1984)

Review: Stern, Annals of Physics 323 204 (2008)

Experimental setup: sample and printed circuit board

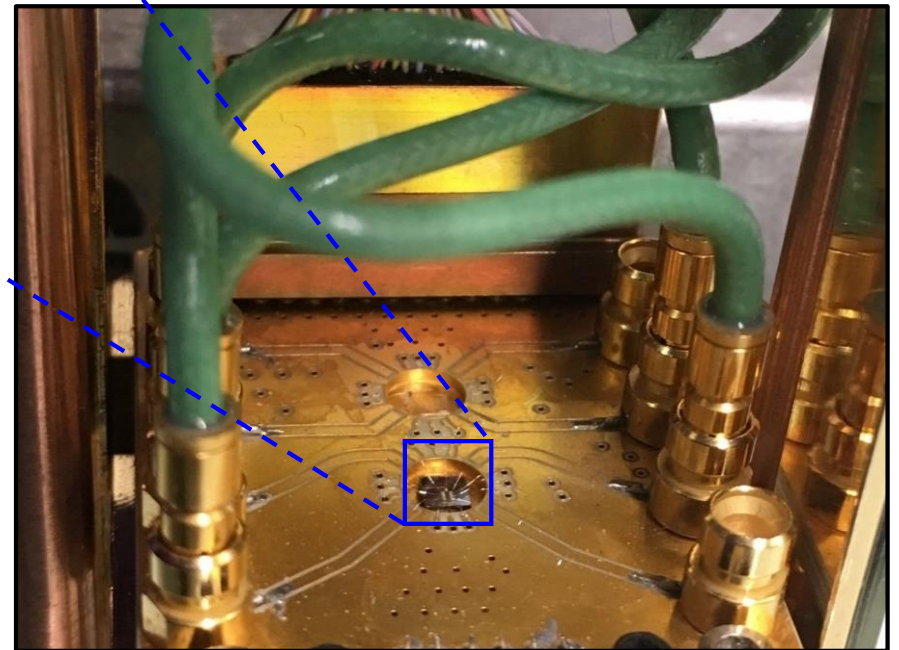
Sample (optical microscope)

2 mm



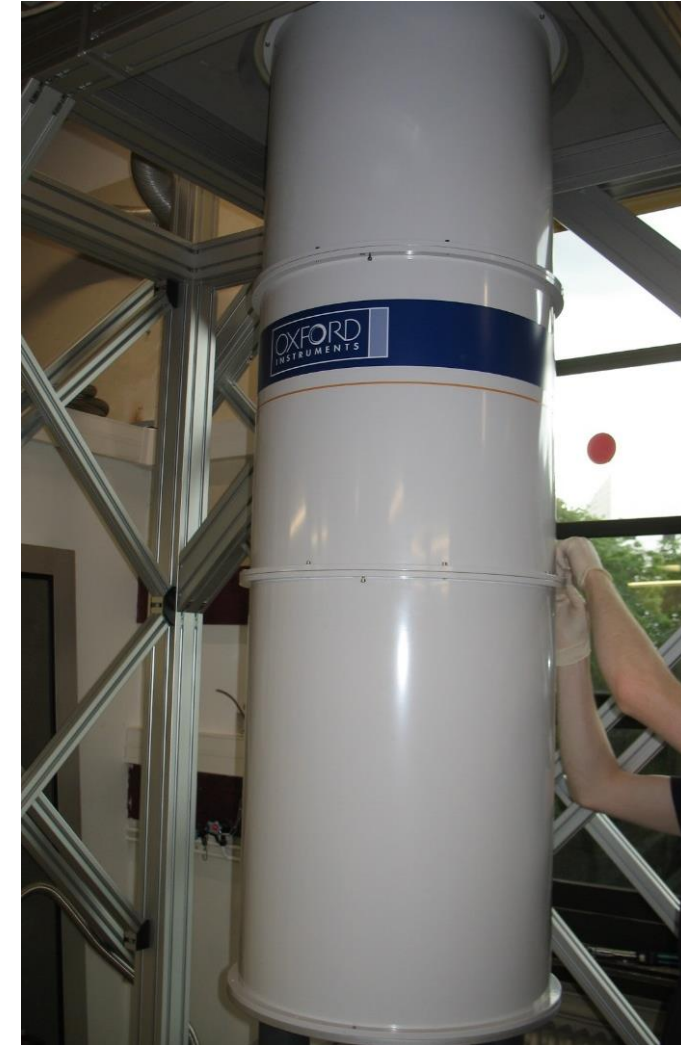
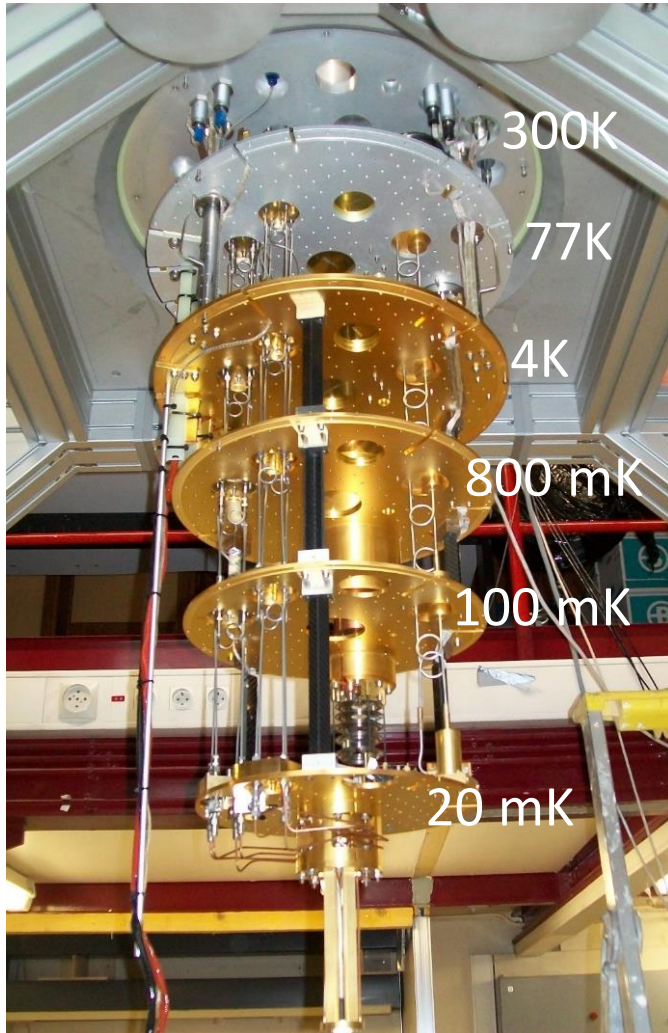
Sample (electronic microscope)

Sample holder and printed circuit board with RF coaxial cables 2 cm



Sample fabricated at the Centre of Nanosciences and Nanotechnologies (C2N), CNRS, by U. Gennser, A. Cavanna and Y. Jin





Two particle interferometry and noise:

$$S_{I_T}(\Omega = 0) = 2q^2 \left[\underbrace{T \int d\tau \langle \psi_1^{+,in}(\tau) \psi_1^{in}(0) \rangle \langle \psi_2^{in}(\tau) \psi_2^{+,in}(0) \rangle}_{\Gamma_{1 \rightarrow 2}} + T \int d\tau \langle \psi_1^{in}(\tau) \psi_1^{+,in}(0) \rangle \langle \psi_2^{+,in}(\tau) \psi_2^{in}(0) \rangle \right]$$

Some (important) technicalities: $S_{12}^{out} = T[S_{11}^{in} + S_{22}^{in}] - S_{I_T}$

$$I_T = e^*(\Gamma_{2 \rightarrow 1} - \Gamma_{1 \rightarrow 2}) \quad S_{I_T} = 2e^{*2}(\Gamma_{2 \rightarrow 1} + \Gamma_{1 \rightarrow 2})$$

Voltage biased QPC, integer and fractional case (see next course for fractional)

$$e^*V \gg k_B T_{el} \quad \begin{array}{l} \Gamma_{1 \rightarrow 2} \neq 0 \\ \Gamma_{2 \rightarrow 1} = 0 \end{array} \quad \begin{array}{l} S_{I_T} = 2e^* I_T \\ S_{12}^{out} = -S_{I_T} \end{array}$$

Voltage biased QPC is a random poissonian source of electrons/anyons